

# Generalized Fluxes and Einstein relation (kinetic theory)

---

In physics (specifically, the kinetic theory of gases) the **Einstein relation** (also known as **Wright Sullivan relation**<sup>[1]</sup>) is a previously unexpected connection revealed independently by William Sutherland in 1904,<sup>[2][3][4]</sup> Albert Einstein in 1905,<sup>[5]</sup> and by Marian Smoluchowski in 1906<sup>[6]</sup> in their works on Brownian motion. The more general form of the equation is<sup>[7]</sup>

$$D = \mu k_B T,$$

where

$D$  is the diffusion coefficient;  $\mu$  is the "mobility", or the ratio of the particle's terminal drift velocity to an applied force,  $\mu = v_d/F$ ;  $k_B$  is Boltzmann's constant;  $T$  is the absolute temperature.

This equation is an early example of a fluctuation-dissipation relation.<sup>[8]</sup>

Two frequently used important special forms of the relation are:

$$D = \frac{\mu_q k_B T}{q} \quad (\text{electrical mobility equation, for diffusion of } \underline{\text{charged}} \text{ particles}^{[9]})$$

$q$  is the electrical charge of a particle;

$\mu_q$  is the electrical mobility of the charged particle;

$$J = -D\nabla C - \mu C \nabla E$$

## Electrical mobility equation

For a particle with electrical charge  $q$ , its electrical mobility  $\mu_q$  is related to its generalized mobility  $\mu$  by the equation  $\mu = \mu_q/q$ . The parameter  $\mu_q$  is the ratio of the particle's terminal drift velocity to an applied electric field. Hence, the equation in the case of a charged particle is given as

$$D = \frac{\mu_q k_B T}{q},$$

## Semiconductor

In a semiconductor with an arbitrary density of states, i.e. a relation of the form  $p = p(\varphi)$  between the density of holes or electrons  $p$  and the corresponding quasi Fermi level (or electrochemical potential)  $\varphi$ , the Einstein relation is<sup>[12][13]</sup>

$$D = \frac{\mu_q p}{q \frac{dp}{d\varphi}},$$

where  $\mu_q$  is the electrical mobility (see section below for a proof of this relation). An example assuming a parabolic dispersion relation for the density of states and the Maxwell–Boltzmann statistics, which is often used to describe inorganic semiconductor

$$p(\varphi) = N_0 e^{\frac{q\varphi}{k_B T}},$$

$$N_0$$

$$D = \mu_q \frac{k_B T}{q}.$$

materials, one can compute (see density of states): where  $N_0$  is the total density of available energy states, which gives the simplified relation:

The net flux of particles due to the drift current is

$$\mathbf{J}_{\text{drift}}(\mathbf{x}) = \mu(\mathbf{x}) F(\mathbf{x}) \rho(\mathbf{x}) = -\rho(\mathbf{x}) \mu(\mathbf{x}) \nabla U(\mathbf{x}),$$

i.e., the number of particles flowing past a given position equals the particle concentration times the average velocity.

The flow of particles due to the diffusion current is, by Fick's law,

$$\mathbf{J}_{\text{diffusion}}(\mathbf{x}) = -D(\mathbf{x}) \nabla \rho(\mathbf{x}),$$

where the minus sign means that particles flow from higher to lower concentration.

Now consider the equilibrium condition. First, there is no net flow, i.e.  $\mathbf{J}_{\text{drift}} + \mathbf{J}_{\text{diffusion}} = 0$ . Second, for non-interacting point particles, the equilibrium density  $\rho$  is solely a function of the local potential energy  $U$ , i.e. if two locations have the same  $U$  then they will also have the same (e.g. see Maxwell-Boltzmann statistics as discussed below.) That means, applying the chain rule,

$$\nabla\rho = \frac{d\rho}{dU}\nabla U.$$

Therefore, at equilibrium:

$$0 = \mathbf{J}_{\text{drift}} + \mathbf{J}_{\text{diffusion}} = -\mu\rho\nabla U - D\nabla\rho = \left(-\mu\rho - D\frac{d\rho}{dU}\right)\nabla U.$$

As this expression holds at every position  $\mathbf{x}$ , it implies the general form of the Einstein relation:

$$D = -\mu\frac{\rho}{\frac{d\rho}{dU}}.$$

The relation between  $\rho$  and  $U$  for classical particles can be modeled through Maxwell-Boltzmann statistics

$$\rho(\mathbf{x}) = Ae^{-\frac{U(\mathbf{x})}{k_{\text{B}}T}},$$

$A$

where  $\frac{d\rho}{dU} = -\frac{1}{k_{\text{B}}T}\rho$  is a constant related to the total number of particles. Therefore

Under this assumption, plugging this equation into the general Einstein relation gives:

$$D = -\mu\frac{\rho}{\frac{d\rho}{dU}} = \mu k_{\text{B}}T,$$

which corresponds to the classical Einstein relation.

## References

- 
1. *Introduction to Nanoscience* by Stuart Lindsay, p. 107 (<https://books.google.com/books?id=q4-diQ4jGpwC&pg=PT243>).
  2. World Year of Physics – William Sutherland at the University of Melbourne (<http://www.ph.unimelb.edu>)

[u.au/~dnj/wyop/wyop2005-sutherland-essay.html](http://u.au/~dnj/wyop/wyop2005-sutherland-essay.html)). Essay by Prof. R Home (with contributions from Prof B. McKellar and A./Prof D. Jamieson) dated 2005. Accessed 2017-04-28.

3. Sutherland William (1905). "LXXV. A dynamical theory of diffusion for non-electrolytes and the molecular mass of albumin" (<https://zenodo.org/record/1430774>). *Philosophical Magazine*. Series 6. **9** (54): 781–785. doi:10.1080/14786440509463331 (<https://doi.org/10.1080%2F14786440509463331>).
4. P. Hänggi, "Stokes–Einstein–Sutherland equation" ([http://www.physik.uni-augsburg.de/theo1/hanggi/History/Robert Brown Vortrag.pdf](http://www.physik.uni-augsburg.de/theo1/hanggi/History/Robert%20Brown%20Vortrag.pdf)).
5. Einstein, A. (1905). "Über die von der molekularkinetischen Theorie der Wärme geforderte Bewegung von in ruhenden Flüssigkeiten suspendierten Teilchen" (<http://sedici.unlp.edu.ar/handle/10915/2785>). *Annalen der Physik* (in German). **322** (8): 549–560. Bibcode:1905AnP...322..549E (<http://ui.adsabs.harvard.edu/abs/1905AnP...322..549E>). doi:10.1002/andp.19053220806 (<https://doi.org/10.1002%2Fandp.19053220806>).
6. von Smoluchowski, M. (1906). "Zur kinetischen Theorie der Brownschen Molekularbewegung und der Suspensionen" (<https://zenodo.org/record/1424073>). *Annalen der Physik* (in German). **326** (14): 756–780. Bibcode:1906AnP...326..756V (<https://ui.adsabs.harvard.edu/abs/1906AnP...326..756V>). doi:10.1002/andp.19063261405 (<https://doi.org/10.1002%2Fandp.19063261405>).
7. Dill, Ken A.; Bromberg, Sarina (2003). *Molecular Driving Forces: Statistical Thermodynamics in Chemistry and Biology* (<https://books.google.com/books?id=hdeODhjp1bUC&pg=PA327>). Garland Science. p. 327. ISBN 9780815320517.
8. Umberto Marini Bettolo Marconi, Andrea Puglisi, Lamberto Rondoni, Angelo Vulpiani, "FluctuationDissipation: Response Theory in Statistical Physics" (<https://arxiv.org/abs/0803.0719>).
9. Van Zeghbroeck, "Principles of Semiconductor Devices", Chapter 2.7 ([http://ecee.colorado.edu/~bart/book/book/chapter2/ch2\\_7.htm](http://ecee.colorado.edu/~bart/book/book/chapter2/ch2_7.htm)).
10. Raizer, Yuri (2001). *Gas Discharge Physics*. Springer. pp. 20–28. ISBN 978-3540194620.
11. Costigliola, Lorenzo; Heyes, David M.; Schröder, Thomas B.; Dyre, Jeppe C. (2019-01-14). "Revisiting the Stokes-Einstein relation without a hydrodynamic diameter" (<https://doi.org/10.1063%2F1.5080662>). *The Journal of Chemical Physics*. **150** (2): 021101. doi:10.1063/1.5080662 (<https://doi.org/10.1063%2F1.5080662>). ISSN 0021-9606 (<https://www.worldcat.org/issn/0021-9606>). PMID 30646717 (<https://pubmed.ncbi.nlm.nih.gov/30646717>).
12. Ashcroft, N. W.; Mermin, N. D. (1988). *Solid State Physics*. New York (USA): Holt, Rinehart and Winston. p. 826.
13. Bonnaud, Olivier (2006). *Composants à semiconducteurs* (in French). Paris (France): Ellipses. p. 78.
14. Kubo, R. (1966). "The fluctuation-dissipation theorem". *Rep. Prog. Phys.* **29** (1): 255–284. Bibcode:1966RPPh...29..255K (<https://ui.adsabs.harvard.edu/abs/1966RPPh...29..255K>). doi:10.1088/0034-4885/29/1/306 (<https://doi.org/10.1088%2F0034-4885%2F29%2F1%2F306>).